

**QUALIFIER  
QUANTUM II**

**Data provided: As needed**

**DEPARTMENT OF PHYSICS NTU**

**February 2025**

**QUANTUM MECHANICS II**

**2 hours**

**Date of Exam:  
Examiner:**

**Answer all questions.  
All question marks summed together, equal 100 points.**

**Question 1. (Time dependent perturbation theory)**

Consider the following Hamiltonian

$$H = H_0 + H'(t) = H_0 + U \exp\left(-\frac{t}{T}\right)$$

where  $U$  is a time independent operator and  $T$  a (decay) constant.

Show that the probability to lowest order in  $U$  that the perturbation will produce a transition from one eigenstate  $n$  of  $H_0$  to a different eigenstate  $m$  of  $H_0$ , during a time interval from  $t = 0$  to a time  $t \gg T$  is given by:

$$|c_m(t)|^2 \xrightarrow{t \gg T} \frac{|\langle m|U|n \rangle|^2}{(E_m - E_n)^2 + \frac{\hbar^2}{T^2}} \quad [15]$$

**Question 2. (perturbation theory)**

A simple one-dimensional harmonic oscillator is subjected to a perturbing linear potential  $\hat{H}_I = b\hat{x}$ , where  $b$  a real number, and  $\hat{x}$  the position.

(a) Calculate the energy shift to the first non-vanishing order of perturbation theory. [10]

(b) Solve the problem exactly and compare with the result obtained in question (a). [10]

Useful:  $\langle n'|\hat{x}|n \rangle = \sqrt{\frac{\hbar}{2m\omega}} (\sqrt{n+1}\delta_{n',n+1} + \sqrt{n}\delta_{n',n-1})$

**Question 3. (Scattering)**

Use the optical theorem to show that there is an upper bound on the elastic scattering cross section  $\sigma_{scattering}$  in the case where the scattering amplitude  $f_k(\theta, \phi)$  is independent of  $\theta$  and  $\phi$ .

(Hint: Optical theorem gives  $\sigma_{scattering} = \frac{4\pi}{k} \text{Im}[f_k(0)]$  ). [15]

**Question 4. (Scattering)**

Start with the equation for  $l$ -scattering cross section

$$\sigma_l = 4\pi |f_l(k)|^2 (2l+1) = 4\pi (2l+1) \left[ \frac{1}{k \cot \delta_l - ik} \right]^2$$

to show that:

$$\sigma_l = \frac{4\pi}{k^2} (2l+1) \frac{(\Gamma/2)^2}{(E - E_R)^2 + (\Gamma/2)^2}$$

where for simplicity we can take  $k = E$ , and  $E_R$  is the resonance energy. Here the resonance width

$\Gamma$ , is defined as follows:  $\frac{1}{\Gamma} = \frac{1}{2} \frac{d\delta_l}{dE}$  [15]

**Question 5. (Scattering)**

In a scattering experiment where a spinless non-relativistic particle of mass  $m$  is scattered by an unknown potential, a resonance is observed where the cross section has a maximum at energy  $E = E_R$ . Use the data of this experiment that are described by the Breit-Wigner formula:

$$\sigma_{scat} = \frac{\pi(2l+1)}{k^2} \frac{\Gamma^2}{(E-E_R)^2 + \frac{\Gamma^2}{4}}$$

to find a value for the orbital angular momentum  $l$  of the resonance. [15]

**Question 6. (Green's functions)**

The unperturbed Schrödinger equation has the form ( $\hbar = 1$ ):

$$\left(i \frac{\partial}{\partial t} + \frac{\nabla^2}{2m}\right) \psi(\vec{x}, t) = 0 \quad (1)$$

with a Kernel function  $K(\vec{x}, t; \vec{y}, t')$  satisfying the homogeneous equation:

$$\left(i \frac{\partial}{\partial t} + \frac{\nabla^2}{2m}\right) K(\vec{x}, t; \vec{y}, t') = 0 \quad (2)$$

with the initial condition:

$$K(\vec{x}, t; \vec{y}, t) = \delta^{(3)}(\vec{x} - \vec{y}) \quad (2a)$$

If we impose the boundary condition that  $K(\vec{x}, t; \vec{y}, t)$  vanishes for  $t' < t$ , we obtain the (*retarded*) Green's function:

$$G_+(\vec{x}, t; \vec{y}, t') = \Theta(t - t') K(\vec{x}, t; \vec{y}, t') \quad (3)$$

satisfying the inhomogeneous equation:

$$\left(i \frac{\partial}{\partial t} + \frac{\nabla^2}{2m}\right) G_+(\vec{x}, t; \vec{y}, t') = i\delta(t - t')\delta^{(3)}(\vec{x} - \vec{y}) \quad (4)$$

Show that the Green's function (3) satisfies the inhomogeneous equation (4). [10]

**Question 7. (Lippmann-Schwinger)**

The “in” states in a scattering problem satisfy the time independent Schrödinger equation:

$$(E - H_0)|\psi_\alpha\rangle^{(+)} = V|\psi_\alpha\rangle^{(+)}$$

show by rewriting the Schrödinger equation above, that a formal solution of the Schrödinger equation is the Lippmann-Schwinger equation:

$$|\psi_\alpha\rangle^{(+)} = |\varphi\rangle + \frac{1}{E - H_0 + i\epsilon} V |\psi_\alpha\rangle^{(+)},$$

where  $|\varphi\rangle$  is the initial state of  $|\psi_\alpha\rangle^{(+)}$  very far from the potential:  $\lim_{t \rightarrow -\infty} |\psi_\alpha\rangle^{(+)} = |\varphi\rangle$  where the potential is negligible.

In this solution we have the presence of an  $i\epsilon$  factor, where  $\epsilon$  a real positive number. Why is this imaginary shift (analytic continuation) in energy introduced? Why is this particular choice used?

Why not  $-i\epsilon$ ? Hint: in your argument you may want to use the time evolution operator

$$U(t, t_0) = \exp(-iE(t - t_0)). \quad [10]$$

**END OF EXAMINATION PAPER**